

## DIFFERENCE SCHEME FOR SEMILINEAR REACTION-DIFFUSION PROBLEM

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### ABSTRACT

We consider the singularly perturbed selfadjoint one-dimensional semilinear reaction-diffusion problem

$$L_\varepsilon y := \varepsilon^2 y''(x) = f(x, y), \quad \text{on } (0,1)$$
$$y(0) = 0; \quad y(1) = 0,$$

where  $f(x,y)$  is a non-linear function. For this problem, using the spline-method with the natural choice of functions, a new difference scheme is given on a non-uniform mesh. The constructed non-linear difference scheme has uniform convergence in points of uneven division segments. A numerical example is given.

**Key words.** Semilinear reaction-diffusion problem, difference scheme, singular perturbation problem

### 1. INTRODUCTION

We consider the semilinear singularly perturbed problem:

$$\varepsilon^2 y''(x) = f(x, y), \quad \text{on } (0,1) \tag{1}$$

$$y(0) = 0; \quad y(1) = 0, \tag{2}$$

where  $0 < \varepsilon < 1$ . Assume that the nonlinear function  $f(x, y)$  is continuously differentiable, and that it has a strictly positive derivative with respect to  $y$ , etc.

$$\frac{\partial f}{\partial y} = f_y \geq m > 0 \quad \text{on } [0,1] \times \mathbb{R} \quad (m = \text{const.}) \tag{3}$$

A solution  $y$  of (1) – (2) usually exhibits sharp boundary layers at the endpoints of  $(0,1)$ , when the parameter  $\varepsilon$  is near zero. When classical numerical methods are applied to (1)-(2), one does not obtain accurate results on the entire interval  $(0,1)$ , because we shall use nonstandard discretisation of (1)-(2).

### 2. CONSTRUCTION OF THE NONLINEAR DIFFERENCE SCHEME

Let us write the differential equation (1) in an equivalent form

$$L_\varepsilon y(x) := \varepsilon^2 y''(x) - \gamma y(x) = \psi(x, y), \quad \text{on } [0,1], \tag{4}$$

where  $\psi(x, y) = f(x, y) - \gamma y$ , and  $\gamma \geq m$  is a chosen constant. Consider an arbitrary grid  $0 = x_0 < x_1 < x_2 < \dots < x_N = 1$ , and consider the following boundary problems

$$L_\varepsilon u_i(x) := 0, \quad \text{on } (x_i, x_{i+1}); \quad u_i(x_i) = 1, \quad u_i(x_{i+1}) = 0, \quad (i = 0, 1, \dots, N-1) \quad (5a)$$

$$L_\varepsilon u_i(x) := 0, \quad \text{on } (x_i, x_{i+1}); \quad u_i(x_i) = 0, \quad u_i(x_{i+1}) = 1, \quad (i = 0, 1, \dots, N-1) \quad (5b)$$

We denote the solutions of problems (5a) and (5b) by  $u_i^I(x)$ ,  $u_i^{II}(x)$  ( $i = 0, 1, \dots, N-1$ ), respectively.

Consider a new boundary problem

$$L_\varepsilon y_i(x) = \psi(x, y_i), \quad \text{on } (x_i, x_{i+1}) \quad (i = 0, 1, \dots, N-1) \quad (6)$$

$$y_i(x_i) = y(x_i) \quad ; \quad y_i(x_{i+1}) = y(x_{i+1}).$$

Clearly, we have  $y_i(x) \equiv y(x)$  on  $[0, 1]$  ( $i = 0, 1, \dots, N-1$ ). The solution of (6) is given by

$$y_i(x) = C_1 u_i^I(x) + C_2 u_i^{II}(x) + \int_{x_i}^{x_{i+1}} G_i(x, s) \psi(s, y(s)) ds, \quad (x \in [x_i, x_{i+1}]),$$

where  $G_i(x, s)$  is the Green's function associated with the operator  $L_\varepsilon$  on segment  $[x_i, x_{i+1}]$ . The function  $G_i(x, s)$  in this case has the following form

$$G_i(x, s) = \frac{1}{\varepsilon^2 w_i(s)} \begin{cases} u_i^{II}(x) u_i^I(s) & ; \quad x_i \leq x \leq s \leq x_{i+1} \\ u_i^I(x) u_i^{II}(s) & ; \quad x_i \leq s \leq x \leq x_{i+1} \end{cases}, \quad (7)$$

where  $w_i(s) = u_i^{II}(s) (u_i^I(x))'_{x=s} - u_i^I(s) (u_i^{II}(x))'_{x=s}$ .

Clearly,  $w_i(s) \neq 0$  ( $s \in [x_i, x_{i+1}]$ ), because solutions  $u_i^I$  and  $u_i^{II}$  are linearly independent.

From boundary conditions in (6), it follows that  $C_1 = y(x_i) = y_i$ ,  $C_2 = y(x_{i+1}) = y_{i+1}$  ( $i = 0, 1, \dots, N-1$ ).

Hence, solution  $y_i$  of (6) on the segment  $[x_i, x_{i+1}]$  has the following form

$$y_i(x) = y_i u_i^I(x) + y_{i+1} u_i^{II}(x) + \int_{x_i}^{x_{i+1}} G(x, s) \psi(s, y(s)) ds. \quad (8)$$

Functions  $u_i^I(x)$  and  $u_i^{II}(x)$  are known from earlier papers (see e.g. [2]), and have forms

$$u_i^I(x) = \frac{\sinh(\beta(x_{i+1} - x))}{\sinh(\beta h_i)}, \quad u_i^{II}(x) = \frac{\sinh(\beta(x - x_i))}{\sinh(\beta h_i)} \quad (x \in [x_i, x_{i+1}]), \quad (9)$$

$$(i = 0, 1, \dots, N-1), \quad \text{where } \beta = \frac{\sqrt{\gamma}}{\varepsilon}, \quad h_i = x_{i+1} - x_i.$$

$$\text{Boundary problem: } L_\varepsilon y(x) := \psi(x, y) \quad \text{on } (0, 1), \quad y(0) = y(1) = 0, \quad (10)$$

has a unique continuously differentiable solution  $y(x)$ . Since  $y_i(x) \equiv y(x)$  on  $[x_i, x_{i+1}]$ , ( $i = 0, 1, \dots, N-1$ ) we have that  $y_i'(x)|_{x=x_i} = y'_{i-1}(x)|_{x=x_i}$ , ( $i = 1, 2, \dots, N-1$ ).

Now, differentiating (8), and also by  $y_i'(x)|_{x=x_i} = y'_{i-1}(x)|_{x=x_i}$ , ( $i = 1, 2, \dots, N-1$ ), we get

$$\begin{aligned} & y_{i-1} (u_{i-1}^I(x))'_{x=x_i} + y_i [ (u_{i-1}^{II}(x))'_{x=x_i} - (u_i^I(x))'_{x=x_i} ] + y_{i+1} [ -(u_i^{II}(x))'_{x=x_i} ] = \\ & = \frac{d}{dx} \left[ \int_{x_i}^{x_{i+1}} G_i(x, s) \psi(s, y(s)) ds - \int_{x_{i-1}}^{x_i} G_{i-1}(x, s) \psi(s, y(s)) ds \right]_{x=x_i}. \end{aligned} \quad (11)$$

where  $y_k = y(x_k)$  ( $k = i-1, i, i+1$ ).

We define  $a_i = -\left(u_{i-1}^I(x)\right)'_{x=x_i}$ ;  $d_i = \left(u_{i-1}^II(x)\right)'_{x=x_i} - \left(u_i^I(x)\right)'_{x=x_i}$ ;  $b_i = -\left(u_i^I(x)\right)'_{x=x_i}$ .

From (9) it follows that:  $a_i = \frac{\beta}{\sinh(\beta h_{i-1})}$ ;  $b_i = \frac{\beta}{\sinh(\beta h_i)}$ ;  $d_i = \frac{\beta}{\tanh(\beta h_{i-1})} + \frac{\beta}{\tanh(\beta h_i)}$ .

Hence, now (11) has the following form:

$$-a_i y_{i-1} + d_i y_i - b_i y_{i+1} = \int_{x_i}^{x_{i+1}} \frac{d}{dx} (G_i(x, s))_{x=x_i} \psi(s, y(s)) ds - \int_{x_{i-1}}^{x_i} \frac{d}{dx} (G_{i-1}(x, s))_{x=x_i} \psi(s, y(s)) ds.$$

After differentiating, the right hand side in the above equation becomes

$$a_i y_{i-1} - d_i y_i + b_i y_{i+1} = \frac{1}{\varepsilon^2} \left[ \int_{x_{i-1}}^{x_i} u_{i-1}^{II}(s) \psi(s, y(s)) ds + \int_{x_i}^{x_{i+1}} u_i^I(s) \psi(s, y(s)) ds \right],$$

$$y_i = 0 \quad ; \quad y_N = 0 \quad \text{za } (i = 0, 1, \dots, N-1). \quad (12)$$

Clearly, we cannot generally explicitly compute the integrals in (12). We approximate the function  $\psi(x, y(x))$ , on the segment  $[x_{i-1}, x_i]$ , by

$$\bar{\psi}_{i-1} = \bar{\psi}(x, y(x)) = \psi\left(\frac{x_{i-1} + x_i}{2}, \frac{\bar{y}_{i-1} + \bar{y}_i}{2}\right) \quad (na [x_{i-1}, x_i]) \quad (i = 1, 2, \dots, N),$$

where  $\bar{y}_i$  ( $i = 1, 2, \dots, N-1$ ) are approximation values of the solution  $y(x)$  of the problem (1) – (3) in points  $x_i$  ( $i = 1, 2, \dots, N-1$ ). Finally from (12) we get the difference scheme

$$a_i \bar{y}_{i-1} - d_i \bar{y}_i + b_i \bar{y}_{i+1} = \frac{1}{\varepsilon^2} \left[ \bar{\psi}_{i-1} \int_{x_{i-1}}^{x_i} u_{i-1}^{II} ds + \bar{\psi}_i \int_{x_i}^{x_{i+1}} u_i^I ds \right] \quad (i = 1, 2, \dots, N-1).$$

From (9), we have

$$\int_{x_{i-1}}^{x_i} u_{i-1}^{II}(s) ds = \frac{1}{\beta} \cdot \frac{\cosh(\beta h_{i-1})}{\sinh(\beta h_{i-1})} - \frac{1}{\beta} \cdot \frac{1}{\sinh(\beta h_{i-1})}, \quad \int_{x_i}^{x_{i+1}} u_i^I(s) ds = \frac{1}{\beta} \cdot \frac{\cosh(\beta h_i)}{\sinh(\beta h_i)} - \frac{1}{\beta} \cdot \frac{1}{\sinh(\beta h_i)}.$$

Hence, our difference scheme has the following form:

$$a_i \bar{y}_{i-1} - d_i \bar{y}_i + b_i \bar{y}_{i+1} = \frac{1}{\gamma} \bar{\psi}_{i-1} (c_i - a_i) + \frac{1}{\gamma} \bar{\psi}_i (c_{i+1} - a_{i+1}), \quad \text{where } c_i = \frac{\beta}{\text{tgh}(\beta h_{i-1})}.$$

We define  $c_i - a_i = \Delta c_i$  i  $c_{i+1} - a_{i+1} = \Delta c_{i+1}$ . We can write the last difference scheme in the form

$$a_i \bar{y}_{i-1} - (c_i + c_{i+1}) \bar{y}_i + a_{i+1} \bar{y}_{i+1} = \frac{1}{\gamma} \bar{\psi}_{i-1} \Delta c_i + \frac{1}{\gamma} \bar{\psi}_i \Delta c_{i+1} \quad (14)$$

where  $b_i = a_{i+1}$  and  $d_i = c_i + c_{i+1}$ . Since  $\psi(x, y) = f(x, y) - \gamma y$ , from (14) we have:

$$a_i \bar{y}_{i-1} - (c_i \bar{y}_i + c_{i+1} \bar{y}_i) + a_{i+1} \bar{y}_{i+1} + \frac{\Delta c_i}{2} (\bar{y}_{i-1} + \bar{y}_i) + \frac{\Delta c_{i+1}}{2} (\bar{y}_i + \bar{y}_{i+1}) = \frac{\Delta c_i}{\gamma} \bar{f}_{i-1} + \frac{\Delta c_{i+1}}{\gamma} \bar{f}_i$$

where  $\bar{f}_{i-1} = f\left(\frac{x_{i-1} + x_i}{2}, \frac{\bar{y}_{i-1} + \bar{y}_i}{2}\right)$ . After some computation, we get:

$$\left(\frac{a_i + c_i}{2}\right) \bar{y}_{i-1} - \left(\frac{c_i + a_i}{2} + \frac{c_{i+1} + a_{i+1}}{2}\right) \bar{y}_i + \left(\frac{c_{i+1} + a_{i+1}}{2}\right) \bar{y}_{i+1} = \frac{\Delta c_i}{\gamma} \bar{f}_{i-1} + \frac{\Delta c_{i+1}}{\gamma} \bar{f}_i$$

If we define  $\frac{a_i + c_i}{2} = r_i$ , then the difference scheme (14) gets the simpler form:

$$r_i \bar{y}_{i-1} - (r_i + r_{i+1}) \bar{y}_i + r_{i+1} \bar{y}_{i+1} = \frac{\Delta c_i}{\gamma} \bar{f}_{i-1} + \frac{\Delta c_{i+1}}{\gamma} \bar{f}_i. \quad (15)$$

### 3. NUMERICAL RESULTS

Consider the following boundary problem

$$\varepsilon^2 y'' = y^3 + y - 10, \quad \text{for } x \in (0,1)$$

$$y(0) = y(1) = 0.$$

We use difference scheme (15) to compute the approximate solution.

Table 1. Error  $E_n$ , and convergence rates  $Ord$  for approximate solution

	$\varepsilon = 2^{-4}$	$\varepsilon = 2^{-7}$	$\varepsilon = 2^{-10}$	$\varepsilon = 2^{-16}$	$\varepsilon = 2^{-22}$	$\varepsilon = 2^{-28}$	$\varepsilon = 2^{-30}$	
$N = 64$	1.54e-3	1.73e-3	1.76e-3	1.76e-3	1.76e-3	1.76e-3	1.76e-3	$E_n$
	2.05	2.01	2.01	2.01	2.01	2.01	2.01	$Ord$
$N = 128$	3.83e-4	4.30e-4	4.36e-4	4.36e-4	4.37e-4	4.37e-4	4.37e-4	$E_n$
	2.01	2.00	2.00	2.00	2.00	2.00	2.00	$Ord$
$N = 256$	9.58e-5	1.07e-4	1.09e-4	1.09e-4	1.09e-4	1.09e-4	1.09e-4	$E_n$
	2.00	2.00	2.00	2.00	2.00	2.00	2.00	$Ord$
$N = 512$	2.39e-5	2.68e-5	2.72e-5	2.72e-5	2.72e-5	2.72e-5	2.72e-5	$E_n$

$E_n = \max_{1 \leq i \leq n} |(y^n)_i - (y^{2n})_i|$ , where is  $(y^{(n)})_i = \bar{y}(x_i)$ , approximate values of the unknown function  $y$ , in  $i$ -th points of mesh,  $n$  is the number of points in the mesh. The convergence rate ( $Ord$ ) is defined by  $Ord = \frac{\ln(E_n) - \ln(E_{2n})}{\ln 2}$ .

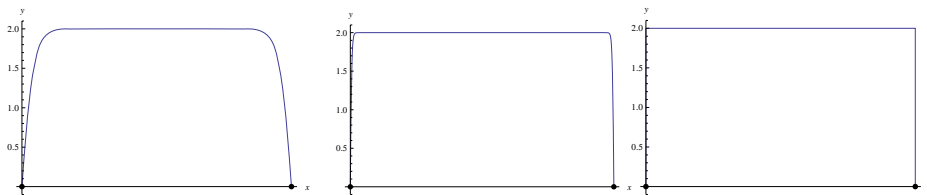


Figure 1. Graphics approximates solutions for values  $\varepsilon = 2^{-4}$ ,  $\varepsilon = 2^{-7}$  and  $\varepsilon = 2^{-30}$

### 4. REFERENCES

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